

# Complex Series

The sum of the first  $N$  terms of a series is

$$S_N = z_1 + z_2 + \dots + z_N = \sum_{n=1}^N z_n$$

Here  $z_n$  can be real, imaginary and complex.

example Maclaurin expansion of  $e^z$

$$S(z) = 1 + z + \frac{1}{2!} z^2 + \frac{1}{3!} z^3 + \dots$$

$$= \sum_{n=0}^{\infty} \frac{z^n}{n!}$$

infinite series

Later, we will learn Taylor series and you will know that  $S(z) = e^z$

In the following, we will introduce

(1) arithmetic series

(2) geometric series

(3) Taylor series — very important 

### arithmetic series

$$a_n = a_0 + nd$$

$$S_N = a_0 + (a_0 + d) + (a_0 + 2d) + \dots + [a_0 + (N-1)d]$$

$$= \sum_{n=0}^{N-1} (a_0 + nd)$$

$$= a_0 \sum_{n=0}^{N-1} 1 + d \sum_{n=0}^{N-1} n$$

$$= a_0 N + d \cdot \frac{1}{2} N(N-1)$$



a piece of cake

### geometric series

$$a_n = a_0 r^n$$

$$S_N = a_0 + a_0 r + a_0 r^2 + \dots + a_0 r^{N-1}$$

$$= \sum_{n=0}^{N-1} a_0 r^n$$

There are many ways to carry out the sum.

$$r^N - 1 = (r-1) (r^{N-1} + r^{N-2} + \dots + r + 1)$$

$$\sum_{n=0}^{N-1} r^n$$

$$\Rightarrow \sum_{n=0}^{N-1} r^n = \frac{r^N - 1}{r - 1} = \frac{1 - r^N}{1 - r}$$

Thus, the sum of a geometric series is

$$S_N = \sum_{n=0}^{N-1} a_0 r^n = a_0 \sum_{n=0}^{N-1} r^n$$

$$= a_0 \frac{1-r^N}{1-r} = a_0 \frac{r^N-1}{r-1}$$

So, the formula seems to carry singularity @  $r=1$ . Is this singularity of  $S_N$  real?

Nope, it is not real ☹

example

$$S = 1 + \frac{2}{2} + \frac{3}{2^2} + \frac{4}{2^3} + \dots$$

Introduce the function  $f(x) = 1 + 2x + 3x^2 + \dots$

It is clear that  $S = f(\frac{1}{2})$ .

Now, we need to carry out the sum for  $f(x)$ .

$$F(x) = \int_0^x f(x') dx' = x + x^2 + x^3 + \dots$$

$$= x(1 + x + x^2 + \dots) = \frac{x}{1-x}$$

Making use of the relation  $f(x) = F'(x)$ ,

$$\begin{aligned}
 f(x) &= \frac{dF}{dx} = \frac{d}{dx} \left[ \frac{x}{1-x} \right] \\
 &= \frac{d}{dx} \left[ -1 + \frac{1}{1-x} \right] \\
 &= \frac{1}{(1-x)^2} \quad \rightarrow \quad S = f\left(\frac{1}{2}\right) = 4
 \end{aligned}$$

$\frac{x-1+1}{1-x} = -1 + \frac{1}{1-x}$

example

$$S(\theta) = 1 + \cos\theta + \frac{\cos 2\theta}{2!} + \dots$$

Making use of the complex algebra  $\ddot{\omega}$

$$S(\theta) = \operatorname{Re} \left\{ 1 + e^{i\theta} + \frac{1}{2!} e^{i2\theta} + \frac{1}{3!} e^{i3\theta} + \dots \right\}$$

Replacing the phase by the complex number,

$e^{i\theta} = z$ , the infinite sum becomes

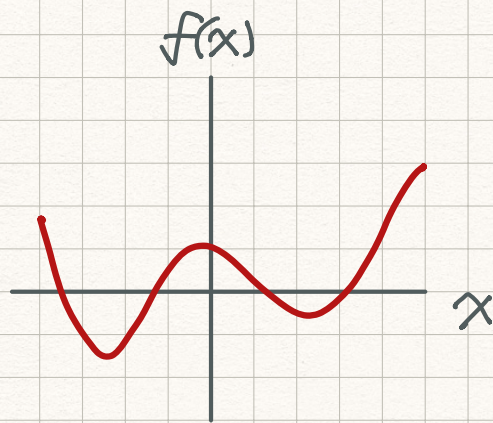
$$S(\theta) = \operatorname{Re} \left\{ 1 + z + \frac{z^2}{2!} + \frac{z^3}{3!} + \dots \right\}$$

$$= \operatorname{Re} \left\{ e^z \right\} = \operatorname{Re} \left\{ e^{\cos\theta + i\sin\theta} \right\}$$

$$= e^{\cos\theta} \cos(\sin\theta) \quad \text{— pretty cool!}$$

## Taylor Series

Given a function  $f(x)$  as shown on the right. Is it possible to express it in terms of power series?



$$f(x) \stackrel{?}{=} a_0 + a_1x + a_2x^2 + a_3x^3 + \dots$$

Skipping the math rigors, one can try to compute the coefficients  $a_n$  ☺

(1) just plug  $x=0$  into  $f(x)$

$$f(0) = a_0 \quad \text{simple...}$$

(2) take a derivative, then plug in  $x=0$

$$f'(0) = a_1 \quad \text{yes!}$$

(3) take two derivatives, then plug in  $x=0$

$$f''(0) = 2a_2 \quad \rightarrow \quad a_2 = \frac{f''(0)}{2!}$$

Keep taking higher-order derivatives,

$$a_n = \frac{f^{(n)}(0)}{n!}$$

Collecting all results together, the fn  $f(x)$  can be expressed as a power series.

$$f(x) = f(0) + f'(0)x + \frac{f''(0)}{2!}x^2 + \dots$$
$$= \sum_{n=0}^{\infty} \frac{f^{(n)}(0)}{n!} x^n$$

This result is remarkable ~ the local properties of the fn [ $f(0), f'(0), f''(0) \dots$  all around  $x=0$ ] dictates its global profile!

example

$$f(x) = e^x$$

$$f'(x) = e^x, \quad f''(x) = e^x, \quad \dots, \quad f^{(n)}(x) = e^x, \quad \dots$$

$$\rightarrow f(0) = f'(0) = f''(0) = \dots = 1$$

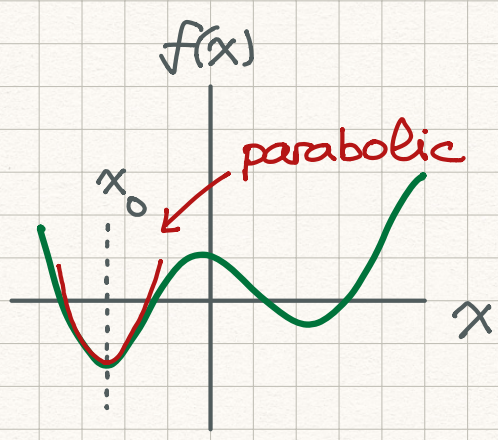
thus,  $e^x = 1 + x + \frac{1}{2!}x^2 + \frac{1}{3!}x^3 + \dots$

One can apply the same trick to expand  $f(x)$  at an arbitrary point  $x = x_0$

$$f(x) = a_0 + a_1(x-x_0) + a_2(x-x_0)^2 + \dots$$

It is straightforward to show that  $a_n = \frac{f^{(n)}(x_0)}{n!}$

$$f(x) = f(x_0) + f'(x_0)(x-x_0) + \frac{f''(x_0)}{2!}(x-x_0)^2 + \dots$$
$$= \sum_{n=0}^{\infty} \frac{f^{(n)}(x_0)}{n!} (x-x_0)^n$$



One can expand  $f(x)$  @ its equilibrium pt [ $f'(x_0) = 0$ ]

According to Taylor series

$$f(x) = f(x_0) + \cancel{f'(x_0)}(x-x_0) + \frac{f''(x_0)}{2!}(x-x_0)^2 + \dots$$

$$\approx f(x_0) + \frac{f''(x_0)}{2!}(x-x_0)^2$$

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quadratic potential  $\ddot{u}$

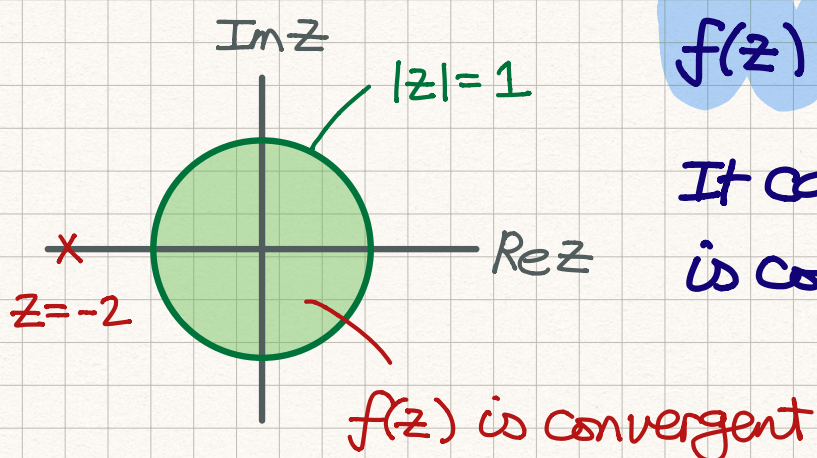
# Analytic Continuation

All the series properties can be generalized to complex numbers on the complex plane.

For instance, consider the infinite sum  $f(z)$

$$f(z) = 1 + z + z^2 + \dots$$

It can be shown that  $f(z)$  is convergent for  $|z| < 1$ .

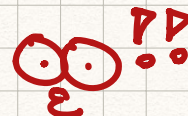


So, does it make any sense when  $z = -2$ ?

$$f(-2) = 1 - 2 + 4 - 8 + \dots$$

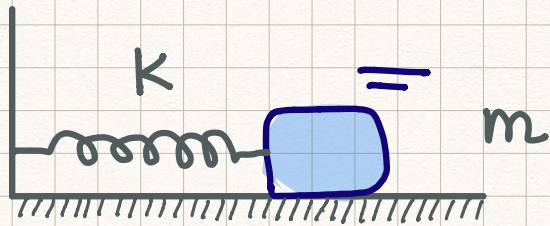
If I tell you that  $1 - 2 + 4 - 8 + \dots = \frac{1}{3}$ ,

maybe you will just drop the applied math course ...





# simple harmonic oscillator



The mass is worn out  
a bit

$$m = m_0 - \Delta m$$

How is the period modified?

Assuming the mass difference is small

$$g = \frac{\Delta m}{m_0} = 1 - \frac{m}{m_0} \ll 1$$

The equation of motion is

$$(m_0 - \Delta m) \ddot{x} + kx = 0$$

Making everything dimensionless  $\ddot{u}$

$$\omega_0^2 = \frac{k}{m_0} \quad \text{so set } \omega_0 t \rightarrow t$$

and introduce  $\alpha = \frac{\omega^2}{\omega_0^2}$

The EOM becomes dimensionless

$$\ddot{x} + x - g\ddot{x} = 0$$

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perturbation!

The solution  $x(t)$  can be expanded,

$$x(t) = x_0(t) + g x_1(t) + g^2 x_2(t) + \dots$$

The zero-th order solution  $x_0(t)$  takes the form

$$x_0(t) = \cos(\sqrt{\alpha} t)$$

$$\alpha = 1 + c_1 g + c_2 g^2 + \dots$$

Let's compute the coefficient  $c_n$ !

①

$$\begin{aligned} (1-g)x'' &= -\alpha(1-g)\cos(\sqrt{\alpha}t) \\ &\quad + (1-g)(g\ddot{x}_1 + g^2\ddot{x}_2 + \dots) \\ &= -\cos(\sqrt{\alpha}t) \\ &\quad + g \left\{ (1-c_1)\cos(\sqrt{\alpha}t) + \ddot{x}_1 \right\} \\ &\quad + g^2 \left\{ (c_1 - c_2)\cos(\sqrt{\alpha}t) + (\ddot{x}_2 - \ddot{x}_1) \right\} \\ &\quad + \dots \end{aligned}$$

②  $x(t) = \cos(\sqrt{\alpha}t) + g x_1 + g^2 x_2 + \dots$

Now, write down EOM order by order  $\ddot{u}$

zero-th order :

trivial.

$$-\cos(\sqrt{\alpha}t) + \cos(\sqrt{\alpha}t) = 0$$

1<sup>st</sup> order :

$$\ddot{x}_1 + x_1 + (1-q)\cos(\sqrt{\alpha}t) = 0$$

$$\rightarrow x_1(t) = \underbrace{A}_{\text{unknown}} \underbrace{\cos(t+\phi)}_{\text{constants}} + \underbrace{B}_{\text{unknown}} \underbrace{\cos(\sqrt{\alpha}t)}_{\text{constants}}$$

substitute back into EOM.

$$(1-\alpha)B \cos(\sqrt{\alpha}t) + (1-q)\cos(\sqrt{\alpha}t) = 0$$

$$\underbrace{(1-\alpha)B}_{\mathcal{O}(g)} + \underbrace{(1-q)}_{\mathcal{O}(1)} = 0$$

$\mathcal{O}(g)$

$\mathcal{O}(1)$

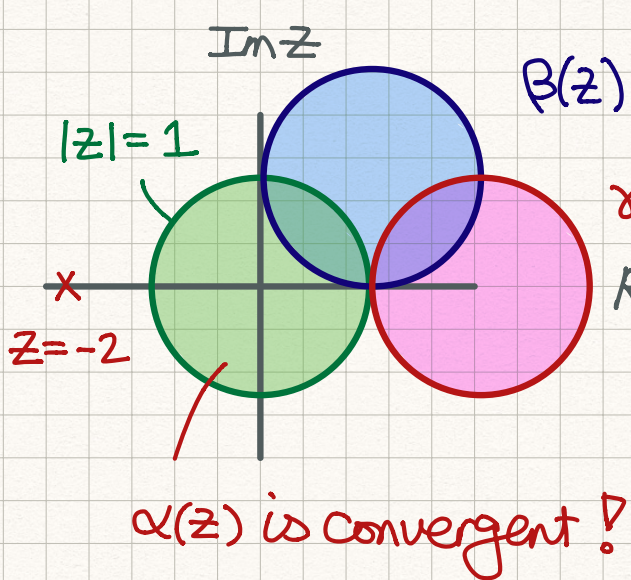
Thus,  $q=1$  and  $B=0$

Because  $x_1(0)=0$  and  $\dot{x}_1(0)=0$ ,

$$A=0 \text{ and } \phi=0 \rightarrow x_1(t)=0$$

One can work out the calculations to all orders.

$$\alpha = \frac{\omega^2}{\omega_0^2} = 1 + g + g^2 + \dots = \sum_{n=0}^{\infty} g^n \quad \leftarrow \frac{1}{1-g}$$



$$\beta(z) = i \{ 1 + i(g-1-i) + i^2(g-1-i)^2 + \dots \}$$

$$\gamma(z) = -1 \{ 1 - (g-2) + (g-2)^2 + \dots \}$$



$$f(z) = \frac{1}{1-z_0} \sum_{n=0}^{\infty} \left( \frac{z-z_0}{1-z_0} \right)^n$$

By analytic continuation on the complex plane,

$$\alpha(z) = \beta(z) = \gamma(z) = \frac{1}{1-z}$$

thus,  $\alpha(z=-2) = 1 - 2 + 4 - 8 + \dots$

$$= \frac{1}{1-(-2)} = \frac{1}{3}$$

Does this make sense?  $m = m_0 - \Delta m = (1-g)m_0$

For  $g = -2$ ,  $m = (1+2)m_0 = 3m_0$

$$\frac{\omega^2}{\omega_0^2} = \frac{k/m}{k/m_0} = \frac{m_0}{m} = \frac{1}{3} \quad \text{YES} \quad \text{🤪}$$