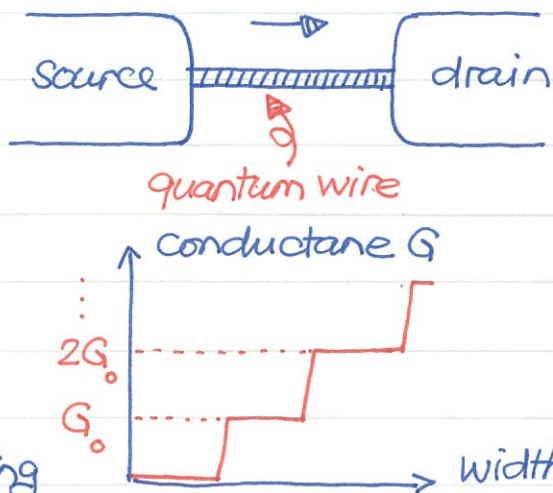


HH0052 Quantum Transport

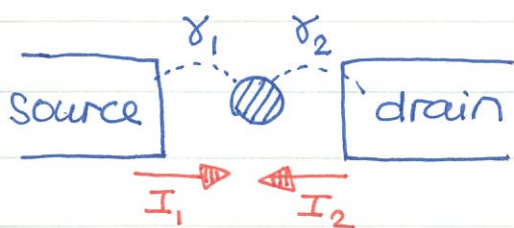
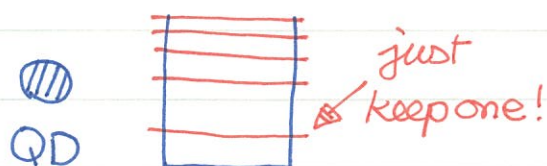
Transport properties are rather different in quantum regime. For instance, the conductance of a clean quantum wire is quantized in units of G_0 ,

$$G_0 = \frac{e^2}{h} = (25.8 \text{ k}\Omega)^{-1}$$

By changing the width of the wire, the conductance increases in steps. How can we understand this surprising phenomena?



① One-level quantum dot: Let us start with a simple system: quantum dot. To simplify the problem, just keep one energy level $E = \epsilon$. The source and drain are described by Fermi-Dirac dist.,



$$f_1(E) = \frac{1}{e^{(E-\mu_1)/\tau} + 1} \quad \text{① } \mu_1 > \mu_2$$

$$f_2(E) = \frac{1}{e^{(E-\mu_2)/\tau} + 1} \quad \text{② } \mu_1 - \mu_2 = qV_{sd}$$



The tunneling rates between the dot and the leads are γ_1 and γ_2 . We can now write down the current Ⓢ

$$I_1 = q\gamma_1 [f_1(1-N) - N(1-f_1)] \rightarrow I_1 = q\gamma_1 [f_1(\epsilon) - N]$$

Similarly, the current from drain into the quantum dot is

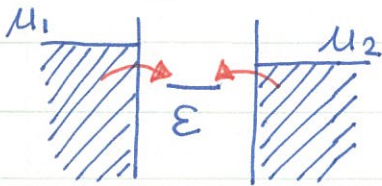
$$I_2 = q\gamma_2 [f_2(1-N) - N(1-f_2)] \rightarrow I_2 = q\gamma_2 [f_2(\epsilon) - N]$$

In steady state, the charge on QD is constant $\rightarrow I_1 + I_2 = 0$

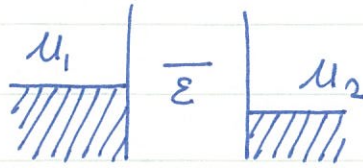
One can then solve for the number N on QD.

$$N = \frac{\gamma_1 f_1 + \gamma_2 f_2}{\gamma_1 + \gamma_2} \rightarrow I = I_1 = -I_2 = q \frac{\gamma_1 \gamma_2}{\gamma_1 + \gamma_2} [f_1(E) - f_2(E)]$$

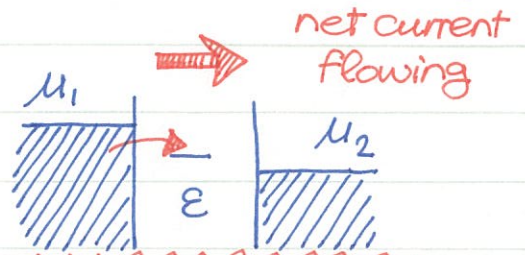
Let's try to understand the tunneling current better. For degenerate Fermi gas, $f_i(E) \approx \Theta(\mu_i - E)$



$f_1(E) = 1, f_2(E) = 1$
no current!



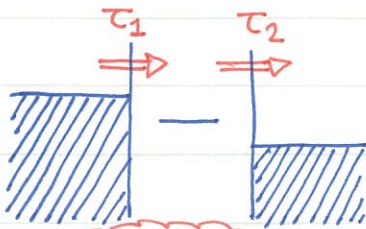
$f_1(E) = 0, f_2(E) = 0$
no current!



$f_1(E) = 1, f_2(E) = 0$

$$I = q \frac{\gamma_1 \gamma_2}{\gamma_1 + \gamma_2} = q \frac{1}{(\frac{1}{\gamma_1}) + (\frac{1}{\gamma_2})}$$

We can understand the current in the case $f_1(E) = 1$ & $f_2(E) = 0$ in the following picture:



$\tau = \tau_1 + \tau_2$

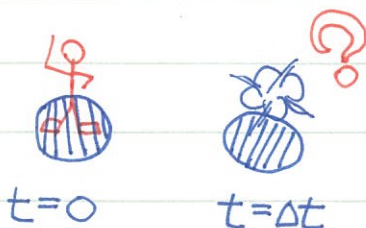
The transport consists two steps:
(Source \rightarrow QD takes time) τ_1
+ (QD \rightarrow drain takes time) τ_2
total time is $\tau = \tau_1 + \tau_2$

$$I = \frac{q}{\tau} = \frac{q}{\tau_1 + \tau_2} = q \frac{1}{\frac{1}{\gamma_1} + \frac{1}{\gamma_2}} = q \frac{\gamma_1 \gamma_2}{\gamma_1 + \gamma_2}$$

Very nice and simple picture ☺

There is only one problem about the above calculations - the result is complete WRONG ☹☹☹

⊗ **Uncertainty principle**: The main problem is the quantum dot is not a closed system. According to uncertainty principle, $\Delta E \Delta t \sim \hbar$



$\rightarrow \Delta E \sim \frac{\hbar}{\Delta t} = \hbar \gamma$ ← uncertainty in energy.

We cannot say the energy is E any more.

Instead, we should use a probability distribution $P(E)$ to describe the single level.

$$\int dE P(E) = 1$$

$P(E) dE =$ probability to find the level in $(E, E+dE)$

After some advanced calculations, the probability density is

$$P(E) = \frac{1}{\pi} \frac{(\hbar\gamma/2)}{(E-\varepsilon)^2 + (\hbar\gamma/2)^2} \xrightarrow{\gamma \rightarrow 0} \delta(E-\varepsilon)$$

Some of you may recognize that the probability density $P(E)$ is just a special case of our good old friend: $D(E)$. Yes, the density of states —

$D(E) dE =$ number of states in $(E, E+dE)$

$$\int dE D(E) = \text{number of ALL states.}$$

It's clear that $P(E) = D(E)$ with (number of all states = 1).

① Re-calculate the tunneling current: Since the energy level is not fixed at $E = \varepsilon$, the tunneling current now becomes

$$I = q \frac{\gamma_1 \gamma_2}{\gamma_1 + \gamma_2} \int dE D(E) [f_1(E) - f_2(E)]$$

Chemical potentials $\mu_1, \mu_2 = \varepsilon \pm \frac{1}{2} q V_{sd}$

For degenerate Fermi gas, $f_i(E) \approx \Theta(\mu_i - E)$

$$\int dE D(E) [f_1(E) - f_2(E)] = \int_{-\infty}^{\mu_1} dE D(E) - \int_{-\infty}^{\mu_2} dE D(E) = \int_{\mu_2}^{\mu_1} dE D(E)$$

$$\rightarrow I = q \frac{\gamma_1 \gamma_2}{\gamma_1 + \gamma_2} \int_{\mu_2}^{\mu_1} dE D(E)$$

The conductance $G = dI/dV_{sd}$ can be computed,

$$\left. \frac{dI}{dV_{sd}} \right|_{V_{sd}=0} = q \frac{\gamma_1 \gamma_2}{\gamma_1 + \gamma_2} \cdot \left[D(\varepsilon) \cdot \frac{1}{2} q - D(\varepsilon) \left(-\frac{1}{2} q\right) \right] = q^2 D(\varepsilon) \frac{\gamma_1 \gamma_2}{\gamma_1 + \gamma_2}$$

Note that the density of states takes the Lorentzian shape and its value at resonance $E = \varepsilon$ is

$$D(\varepsilon) = \frac{1}{\pi} \frac{(\hbar\gamma/2)}{(\varepsilon-\varepsilon)^2 + (\hbar\gamma/2)^2} = \frac{2}{\pi\hbar\gamma} = \frac{4}{h(\gamma_1 + \gamma_2)}$$

$\gamma = \gamma_1 + \gamma_2$

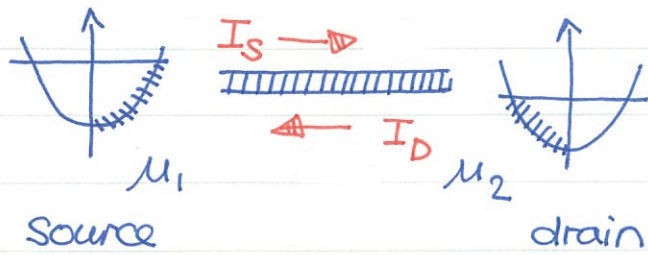
Substitute into the expression for the conductance,

$$G = q^2 D(\epsilon) \frac{\gamma_1 \gamma_2}{\gamma_1 + \gamma_2} \rightarrow \boxed{G = \frac{q^2}{h} \frac{4\gamma_1 \gamma_2}{(\gamma_1 + \gamma_2)^2} \leq \frac{q^2}{h}}$$

The maximum conductance occurs at the symmetrical point $\gamma_1 = \gamma_2$,

$G_{\max} = G_0 = q^2/h$! \leftarrow The conductance has an upper limit (G_0) and does not go to infinity!

① 1D wire : Now we are ready to compute the 1D wire case.



$$I_S = q \int_0^\infty \frac{dp}{2\pi\hbar} f_1(E) \cdot v$$

$$= \frac{q}{2\pi\hbar} \int_0^\infty dE f_1(E) \cdot \left(\frac{dp}{dE} \cdot v \right) \quad \leftarrow I!$$

Similarly, $I_D = q \int_{-\infty}^0 \frac{dp}{2\pi\hbar} f_2(E) \cdot v = - \frac{q}{2\pi\hbar} \int_0^\infty dE f_2(E)$

Total current $I = I_S + I_D = \frac{q}{2\pi\hbar} \int_0^\infty dE (f_1 - f_2) = \frac{q}{2\pi\hbar} (\mu_1 - \mu_2)$

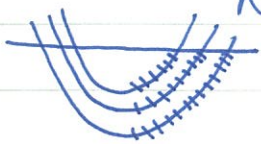
Making use of the relation $\mu_1 - \mu_2 = q V_{sd}$,

$$I = \frac{q^2}{2\pi\hbar} V_{sd} = \frac{q^2}{h} V_{sd} \rightarrow \boxed{G = G_0 = \frac{q^2}{h}}$$

\leftarrow One conducting channel, one $G_0 = q^2/h$!

Generalization to multiple bands is trivial — just count bands.

$N=3$ (number of active bands) $I = N \cdot (q^2/h) \cdot V_{sd}$.



\rightarrow $G = N G_0$ explains the conductance steps.

Note that no impurity scattering is included in the above at all.



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